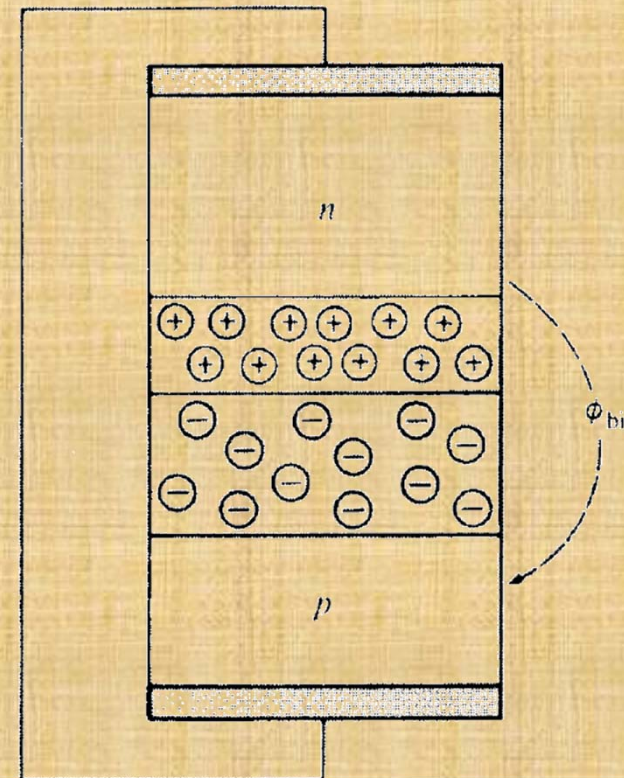


The pn junction

The pn junction

- Abrupt and Graded junction
- Zero bias condition
 - Zero external bias
 - Free electron and hole movement
 - Charged ions
 - Depletion region approximation
 - Electric field formation – Anti-movement mechanism
 - Zero Current
 - Three contact potential
 - Zero electrostatic potential across the device!



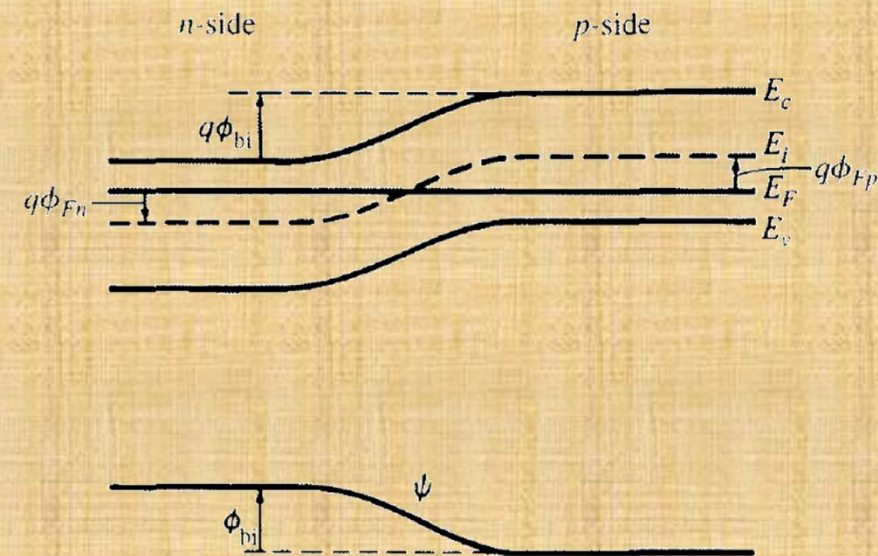
The pn junction

- Zero bias contact potential—Built in potential

$$\phi_{bi} = \phi_{Fp} - \phi_{Fn}$$

- Assume that the n-side is degenerate, and the doping of the p-side is 10^{17} cm^{-3} . Find the value of the contact potential ϕ_{bi} at room temperature:

$$\phi_{bi} = \phi_{Fp} - \phi_{Fn} = 0.42 - (-0.56) = 0.98$$

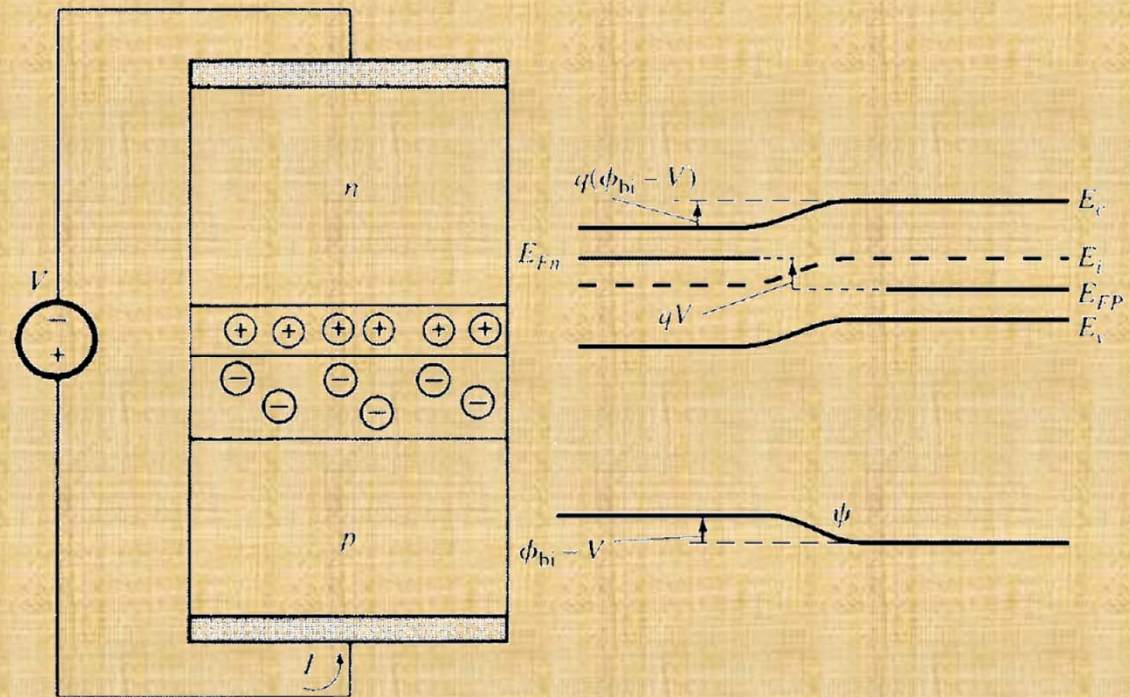


The pn junction...

- Forward bias: $V > 0$
- Equilibrium is now destroyed
- External field is opposite the internal field.
- Current established due to decrease of potential barrier
- How can current flow in depletion region?
- Why it increases exponentially?

$$I = I_0(e^{(V/\phi_T)} - 1)$$

- I_0 is a quantity dependent on junction geometry and physical parameters of the semiconductor material and is an increasing function of temperature
- Depletion region reduction



The pn junction...

- Reverse bias : $V < 0$
- Same as what happen in MOS transistor
- Increase of electrostatic potential across the depletion region:

$$\phi_c = \phi_{bi} + V_R$$

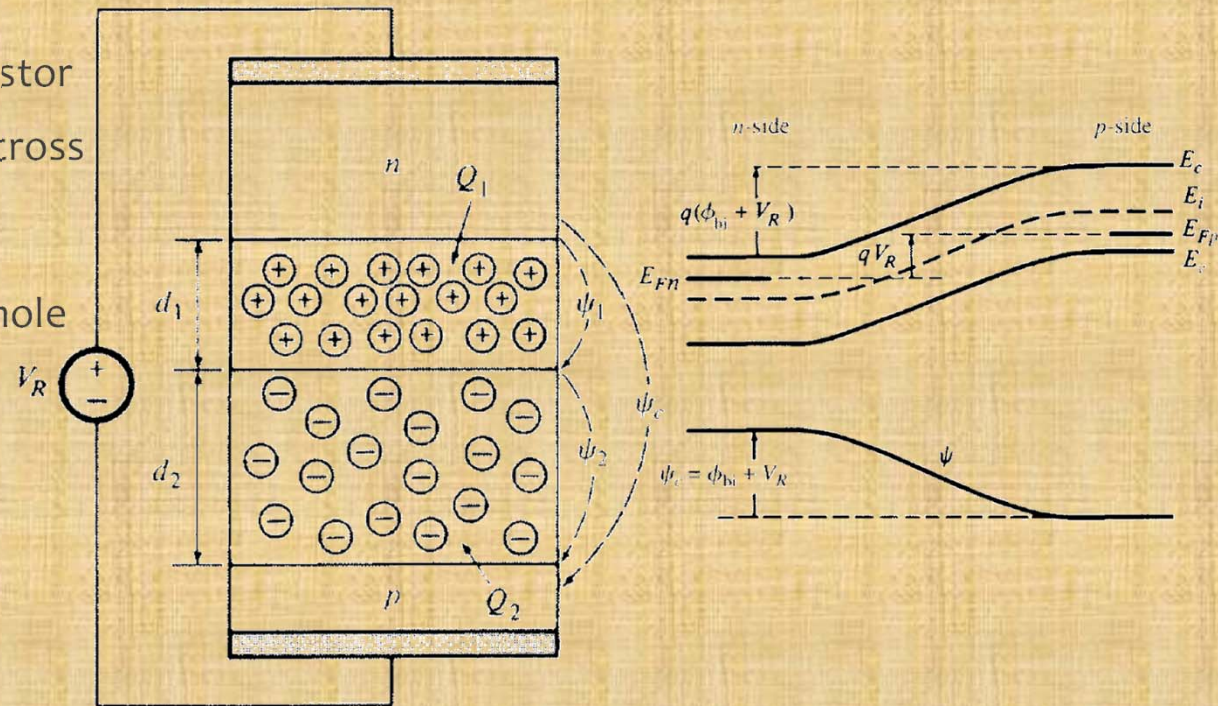
- More band bending—electron and hole movement

- Local charge densities:

$$Q_1 = +q(d_1 A)N_D$$

$$Q_2 = -q(d_2 A)N_A$$

$$Q_1 = -Q_2 \Rightarrow d_1/d_2 = N_A/N_D$$



The pn junction...

- Calculation of electric field in the depletion region

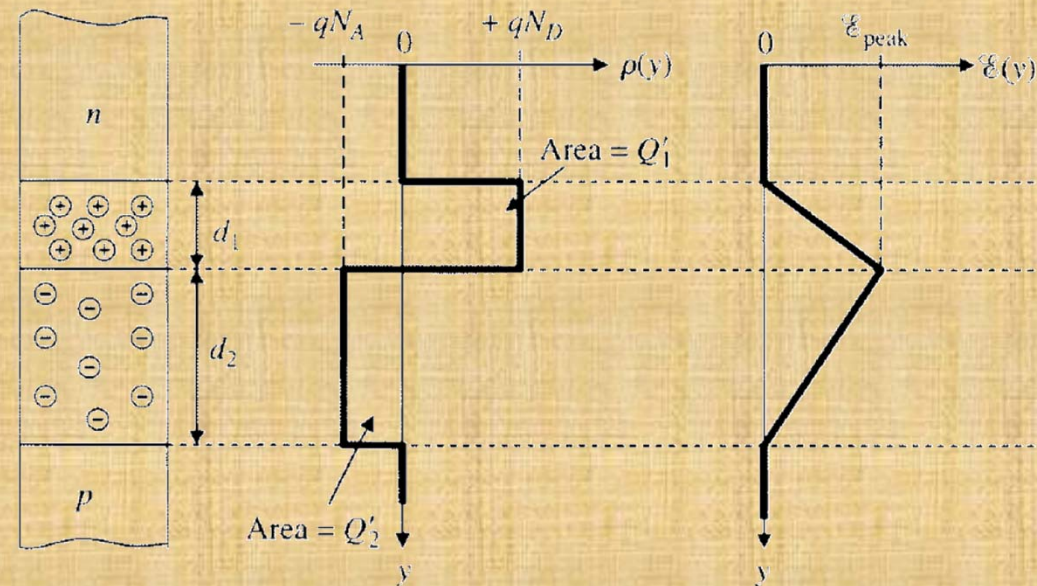
$$\frac{dE(x)}{dx} = \frac{\rho}{\epsilon_s}$$

$$E(y) = E(y_0) + \frac{1}{\epsilon_s} \int_{y_0} \rho(y') dy'$$

$$\rho = qN_D \text{ for } Q'_1$$

$$E_{\max(\text{peak})} = \frac{qN_D d_1}{\epsilon_s}$$

$$= \frac{qN_A d_2}{\epsilon_s}$$



The pn junction...

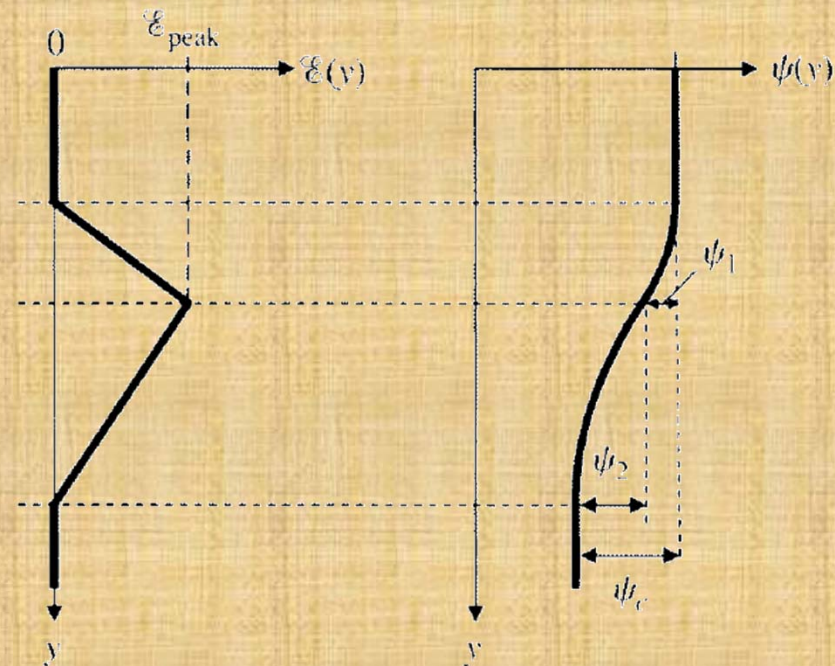
- Calculation of electrostatic potential in the depletion region:

$$\psi(y) = \psi(y_0) - \int_{y_0}^y E(y') dy'$$

$$\psi_1 = \frac{E_{max} d_1}{2} = \frac{q N_D d_1^2}{2 \epsilon_s}$$

$$\psi_2 = \frac{E_{max} d_2}{2} = \frac{q N_A d_2^2}{2 \epsilon_s}$$

$$\psi_c = \psi_1 + \psi_2$$



The pn junction...

- Equation for highly doped sided such as n+p junction we have:

$$N_D \gg N_A \Rightarrow d_1 \ll d_2$$

$$\psi_1 \ll \psi_2 \Rightarrow \psi_c = \psi_2$$

- Nearly all voltage drop will be on **low level doping** region

$$d_2 = \sqrt{\frac{2\epsilon_s}{qN_A}} \sqrt{\psi_c}$$
$$Q_2' = \frac{Q_2}{A} = \sqrt{2q\epsilon_s N_A} \sqrt{\psi_c}$$

- Minority and Majority carrier motion! – Small current flow. Why?
- Saturation current – $V < 3\Phi_T$

$$I = -I_0$$

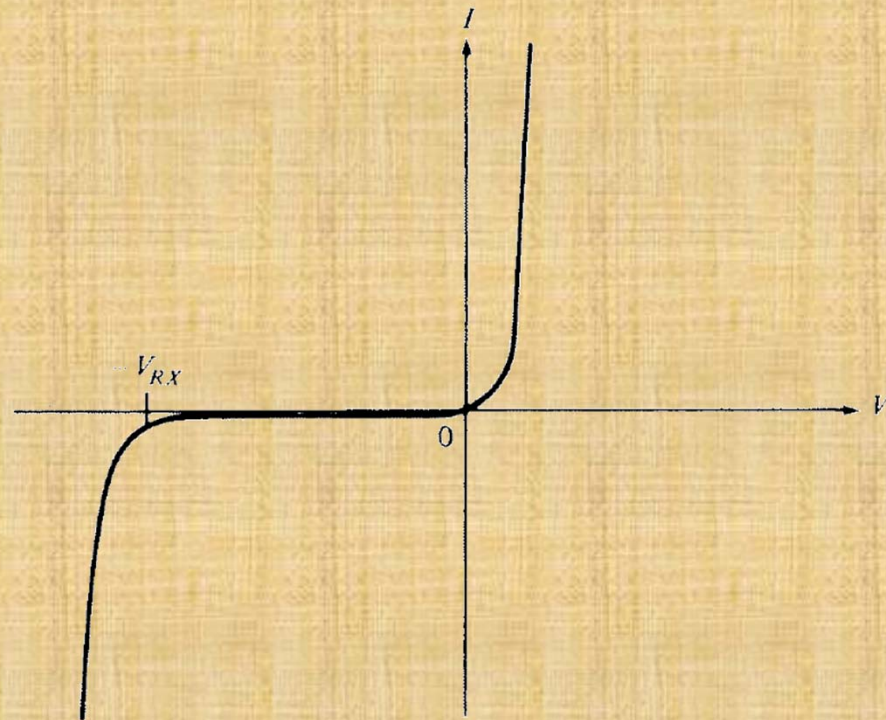
- Can we assume it as a constant value?
- Saturation current dependency to temperature -- double for every 8°C
- Reverse break down voltage– V_{RX} from 5 V to 100 V



The pn junction...

- I/V characteristics of pn junction

$$I = I_0(e^{(V/\phi_T)} - 1)$$



The pn junction...

- Junction capacitance – Small signal overview
- The change in reverse voltage \Rightarrow Change in depletion region
- If ΔV_R Increases then depletion region increases and vice versa
- As we know: $C = \frac{Q}{V}$
- For small signal capacitance we have:

$$C_j = \frac{\Delta Q}{\Delta V_R} = \frac{dQ}{dV_R}$$
$$C_j' = \frac{dQ'}{dV_R} = \frac{\sqrt{2q\epsilon_s N_A}}{2\sqrt{V_R + \phi_{bi}}} \text{ or } \frac{\epsilon_s}{d_2}$$
$$C_j' = \frac{C_{j0}'}{\sqrt{\frac{V_R}{\phi_{bi}} + 1}}$$



The pn junction...

- Other type of junction capacitance:

$$C_j' = \frac{C_{j0}'}{\left(\frac{V_R}{\phi_{bi}} + 1\right)^\alpha}$$

- α is $\frac{1}{2}$ for abrupt highly doped junction
- α is $\frac{1}{3}$ for gradually linearly doped junction
- for other type of junction can be extracted from fitting results

C_j' (F/cm²)

